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$$(a) \hat{A} = \frac{d}{du} \quad \hat{E} = u$$

$$\begin{aligned} \hat{A}^2 f(u) &= \hat{A}[\hat{A}f(u)] = \hat{A}\left[\frac{d}{du}f(u)\right] \\ &= \frac{d}{du}\left[\frac{d}{du}f(u)\right] = \frac{d^2}{du^2}f(u) \end{aligned}$$

$$\hat{A}^2 \equiv \frac{d^2}{du^2}$$

$$[\hat{A}f(u)]^2 = \left[\frac{d}{du}f(u)\right]^2 \neq \hat{A}^2 f(u)$$

$$\begin{aligned} (b) \hat{A}\hat{E}f(u) &= \frac{d}{du}[uf(u)] \\ &= uf(u) + u^2 \frac{df(u)}{du} \end{aligned}$$

$$\begin{aligned} \hat{E}\hat{A}f(u) &= u^2 \left[\frac{d}{du}f(u)\right] \\ &= u^2 \frac{df(u)}{du} \end{aligned}$$

$$\text{hence } \hat{A}\hat{E}f(u) \neq \hat{E}\hat{A}f(u)$$

7.

(a)  $e^{iku} \Rightarrow$  eigenfns of  $\frac{d}{du}$  and  $\frac{d^2}{du^2}$

$$\frac{d}{du} e^{iku} = \underbrace{(ik)}_{\text{eigenvalue}} e^{iku}$$

$$\frac{d^2}{du^2} e^{iku} = \underbrace{(ik)^2}_{\text{eigenvalue}} e^{iku} = \underbrace{-k^2}_{\text{eigenvalue}} e^{iku}$$

(b)  $\frac{d}{du} \cos ku = -k \sin ku$   
 $\sin ku$  is not the original function  
 $\cos ku$  is not an eigenfunction of  $\frac{d}{du}$

$$\frac{d^2}{du^2} \cos ku = \underbrace{-k^2}_{\text{eigenvalue}} \underbrace{\cos ku}_{\text{original fn. comes back}}$$

(c)  $\frac{d}{du} k = 0 = \underbrace{0}_{\text{eigenvalue}}, k$

$$\frac{d^2}{du^2} k = 0 \quad \text{eigenvalue} = 0$$

7 (d)  $\frac{d}{du} (ku) = k = \frac{1}{k} (ku)$

$\uparrow$  original fn.  
 $\uparrow$  not a constant; hence cannot be an eigenvalue

$ku$  is not an eigenfn. of  $\frac{d}{du}$

$\frac{d^2}{du^2} (ku) = 0 = 0 \cdot (ku)$

$\uparrow$  eigenvalue

$ku$  is an eigenfn. of  $\frac{d^2}{du^2}$

(e)  $e^{-au^2} = f(u)$

$\frac{d}{du} (e^{-au^2}) = -2au (e^{-au^2})$

$\underbrace{\hspace{10em}}$  not an eigenfn. of  $\frac{d}{du}$   
 $\uparrow$  not a constant; hence not an eigenvalue

$\frac{d^2}{du^2} (e^{-au^2}) = \frac{d}{du} \{ -2au (e^{-au^2}) \}$

$= (-2a + 4a^2 u^2) e^{-au^2}$

$\underbrace{\hspace{10em}}$  not an eigenvalue

$e^{-au^2}$  is not an eigenfn. of  $\frac{d^2}{du^2}$

118. 

$$\dagger \quad \frac{d^2}{du^2} \cos \omega u = \underline{\underline{-\omega^2 \cos \omega u}}$$

$$\dagger \quad \frac{d}{dt} e^{i\omega t} = \underline{\underline{i\omega e^{i\omega t}}}$$

$$\dagger \quad \left( \frac{d^2}{du^2} + 2 \frac{d}{du} + 3 \right) e^{\alpha u}$$

$$= \underline{\underline{(d^2 + 2\alpha + 3) e^{\alpha u}}}$$

$$\dagger \quad \frac{\partial}{\partial \gamma} x^2 \exp(6\gamma) = \underline{\underline{6 x^2 \exp(6\gamma)}}$$

9. 

$$\nabla^2 (\cos ax) (\cos by) (\cos cz)$$

$$= \left[ \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right] \{ (\cos ax) (\cos by) (\cos cz) \}$$

$$= -a^2 (\cos ax) (\cos by) (\cos cz) - b^2 (\cos ax) (\cos by) (\cos cz)$$

$$- c^2 (\cos ax) (\cos by) (\cos cz)$$

$$= \underline{\underline{- (a^2 + b^2 + c^2) (\cos ax) (\cos by) (\cos cz)}}$$

↑  
eigen value

15.

(a)  $\int_0^L N \sin \frac{n\pi x}{L} N \sin \frac{n\pi x}{L} dx = 1$  wave fun. for a particle in a 1-D box

$$N^2 \int_0^L \sin^2 \frac{n\pi x}{L} dx = 1$$

[using ~~cos 2x = 1 - 2sin^2 x~~  $\cos 2x = 1 - 2\sin^2 x$

$$\therefore \sin^2 x = \frac{1}{2} (1 - \cos 2x)]$$

$$\therefore \frac{N^2}{2} \int_0^L [1 - \cos \frac{2n\pi x}{L}] dx = 1$$

$$\frac{N^2}{2} \left[ \int_0^L dx - \int_0^L \cos \frac{2n\pi x}{L} dx \right] = 1$$

$$\frac{N^2}{2} \left[ L - \left\{ \frac{L}{2n\pi} \sin \frac{2n\pi x}{L} \right\}_0^L \right] = 1$$

$$\frac{N^2}{2} \left[ L - \frac{L}{2n\pi} \underbrace{\left\{ \sin 2n\pi \right\}}_0 \right] = 1$$

$$\therefore \frac{N^2}{2} L = 1$$

$$\therefore N = \sqrt{\frac{2}{L}}$$

(b)

$$N^2 \int_{-L}^L c^2 dx = 1$$

$$N^2 c^2 \int_{-L}^L dx = 1$$

$$\therefore 2LN^2 c^2 = 1$$

$$\therefore N = \frac{1}{c\sqrt{2L}}$$

(c)

Here you will have to use spherical polar coordinates.

$$\frac{dx dy dz}{\downarrow \text{Cartesian coordinates}} = d\tau = \underbrace{r^2 \sin\theta dr d\theta d\phi}_{\downarrow \text{spherical polar coordinates}}$$

$$N^2 \int_0^{\infty} r^2 e^{-r/a_0} \cdot e^{-r/a_0} dr \int_0^{\pi} \sin\theta d\theta \int_0^{2\pi} d\phi = 1$$

$$\therefore N^2 \int_0^{\infty} r^2 e^{-2r/a_0} dr \cdot (2) \cdot (2\pi) = 1$$

$$\therefore 4\pi N^2 \int_0^{\infty} r^2 e^{-2r/a_0} dr = 1$$

Standard integral given with no question

$$\left[ \int_0^{\infty} x^n e^{-ax} dx = \frac{n!}{a^{n+1}} \right]$$

$$\therefore 4\pi N^2 \left( \frac{a_0^3}{4} \right) = 1 \quad \therefore N = \frac{1}{(\pi a_0^3)^{1/2}}$$

(d)

$\kappa \exp(-r/2a_0)$  in 3D

$$\underline{\kappa = r \sin \theta \cos \phi}$$

$$\therefore N^2 \psi^2 = N^2 r^2 \sin^2 \theta \cos^2 \phi e^{-r/a_0}$$

$$\therefore N^2 \int_0^{\infty} r^2 \cdot r^2 e^{-r/a_0} dr \int_0^{\pi} \sin^3 \theta d\theta \int_0^{2\pi} \cos^2 \phi d\phi = 1$$

Do this integral by yourself

Answer:

$$N = \frac{1}{(32\pi a_0^5)^{1/2}}$$

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17. (a)  $e^{-x^2/2}$   $(-\infty, \infty)$

both at  $-\infty$  and  $+\infty$   $e^{-x^2/2}$  converge; hence  
psi is a valid wavefn and so is normalizable.

To normalize

$$\int \psi^* \psi dx = 1$$

$$\text{or } \int_{-\infty}^{+\infty} N^* e^{-x^2/2} N e^{-x^2/2} dx = 1$$

$$\text{or } |N|^2 \int_{-\infty}^{+\infty} e^{-x^2/2} dx = 1$$

$$\text{or } 2|N|^2 \int_0^{\infty} e^{-x^2/2} dx = 1$$

└──────────┘  
Standard integral

[ $e^{-x^2/2}$  is an  
odd function]

$$\text{or } 2|N|^2 \left(\frac{\pi}{4}\right)^{1/2} = 1$$

$$\therefore |N|^2 \cdot \pi^{1/2} = 1$$

$$\therefore |N|^2 = \frac{1}{\pi^{1/2}}$$

$$\therefore N = \frac{1}{\pi^{1/4}}$$

[neglecting -ve sign]

17 (c)  $e^{i\theta}$  ( $0, 2\pi$ )  $\rightarrow$  can be normalized

$$N^2 \int_0^{2\pi} e^{-i\theta} e^{i\theta} d\theta = 1$$

$$\therefore N = \left(\frac{1}{2\pi}\right)^{1/2}$$

$$\psi(e) N^2 \int_0^\infty dn \int_0^\infty dy e^{-(n+iy)^2/2} e^{-(n-iy)^2/2}$$

$$\therefore N^2 \int_0^\infty e^{-n^2} dn \int_0^\infty e^{-y^2} dy = 1$$

$$\text{or } N^2 \left(\frac{\pi}{4}\right)^{1/2} \left(\frac{\pi}{4}\right)^{1/2} = 1$$

$$\therefore N^2 = \frac{4}{\pi}$$

$$\text{and } \underline{\underline{N = \frac{2}{\sqrt{\pi}}}}$$

wavefn. converges at the limits of 0 and  $\infty$ , hence

it could be normalized

(b)  $e^x$   $(0, \infty)$

the function diverges at  $x \rightarrow \infty$

hence  $e^x$  cannot be normalized

(d)  $x e^x$   $(0, \infty)$

here also the function diverges  $\rightarrow$  cannot be normalized

however if the function had been

$$x e^{-x} (0, \infty)$$

then it would have been normalizable

18.

$$\psi(x) = \frac{1}{2} \phi_1(x) + \frac{1}{4} \phi_2(x) + \frac{3+i\sqrt{2}}{4} \phi_3(x)$$

$\phi_1(x)$ ,  $\phi_2(x)$  and  $\phi_3(x)$  are normalized

given  $\hat{H} \phi_1(x) = E_1 \phi_1(x)$   
 considering  $\hat{H}$  to be the K.E. operator  $\hat{H} \phi_2(x) = 3E_1 \phi_2(x)$   
 $\hat{H} \phi_3(x) = 7E_1 \phi_3(x)$

(b) these are the 3 possible values for K.E.

(a)  $\int_{\text{all space}} \psi^* \psi dx = |c_1|^2 + |c_2|^2 + |c_3|^2$

$$= \frac{1}{4} + \frac{1}{16} + \frac{11}{16} = \frac{16}{16} = 1$$

Hence  $\psi(x)$  is also normalized

(d)  $\langle E \rangle = \frac{\int \psi^* \hat{H} \psi dx}{\int \psi^* \psi dx} = 1$  [found above]

$$= \int \psi^* \hat{H} \psi$$

$$= |c_1|^2 E_1 + |c_2|^2 3E_1 + |c_3|^2 7E_1$$

↑  
↑  
↑  
respective eigen values

$$= \frac{E_1}{4} + \frac{3E_1}{16} + \frac{11}{16} \cdot 7E_1$$

$$= \frac{84E_1}{16} = \underline{\underline{5.25E_1}}$$

(e)  $P = |c_k|^2$  i.e.  $\frac{1}{4}$  for  $(E_1)$ ;  $\frac{1}{16}$  for  $(3E_1)$ ;  $\frac{11}{16}$  for  $(7E_1)$